#### Swampland, Stringy de Sitter, and the Measure Problem

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with focus on recent work with Friedrich / Walcher / Westphal

<u>Outline</u>

- The flux landscape, the Swampland, and Stringy de Sitter
- Problems of KKLT/LVS. Can stringy quintessence help?
- The measure problem: a re-evaluation in view of 'Rocky' and 'Swampy' landscapes
- The (potentially) key role of the Cobordism Conjecture and End-of-the-Worlds branes

• Towards predictions....

String Compactifications

• String theory provides an (essentially unique) and UV-complete field theory in 10d:

$$S_{IIB} = \int_{10} \mathcal{R} - |F_{\mu\nu\rho}|^2 + \cdots$$

[Let's say type IIB, to be concrete.]

- Compactifying on Calabi-Yau-Orientifolds, one preserves  $\mathcal{N} = 1$  SUSY and (classically) zero 4d cosmological constant.
- The extra ingredient of fluxes induces an exponentially large landscape of discrete solutions.



Bousso/Polchinski, Giddings/Kachru/Polchinski, Denef/Douglas '04

### String compactifications: flux landscape

 One usually visualizes the emerging situation as follows: (just with φ → {φ<sub>1</sub>, · · · , φ<sub>N</sub>})



- But this picture jumps very far ahead.
- So far we only stabilized the shape ('complex structure') moduli.
- Classically, the size ('Kahler') moduli remain flat and the CC of all vacua is zero.

#### String compactifications: beyond leading order

- The size moduli (let's say just the volume) get a (much smaller) potential from quantum corrections.
- All known effects are of exponential runaway type.
- Two such effects can give SUSY-AdS.
- It takes a conspiracy between at least three 'runaway potentials' to get meta-stable de-Sitter vacua.



The historical prime example: KKLT

Kachru/Kallosh/Linde/Trivedi

- Assume there is just one of flat direction: the volume.
- Its potential comes from a non-perturbative effect:



 $\Rightarrow$   $W = W_0 + e^{-T}$ , (where  $W_0$  is the previous flux effect)

- $\Rightarrow$  V  $\sim e^{-2T} |W_0|e^{-T}$ 
  - ⇒ Kahler modulus stabilized (controlled for  $W_0 \ll 1$ ).

Ads minimum

• This AdS model is 'Step 1' of KKLT.

# KKLT (continued)

- 'Step 2' involves 'uplifting' to dS by adding an anti-D3-brane.
- This requires a 'strongly warped region' or 'Klebanov-Strassler throat' (met audi (realized by introducing a large amount of flux in an appropriate (conifold) region of the CY).



• Eventually, one may hope for the desired potential:



#### But:

Full explicitness has remained elusive for technical reasons.

#### The swampland (counter?) revolution and the dS conjecture

- This, and some important variants (like 'LVS') has remained the main evidence for 'stringy dS'.
- No analogues in type-I, IIA, heterotic, 11d SUGRA were found.
- Based on this, it has been proposed that stringy dS is impossible as a matter of principle ('is in the Swampland').

Danielsson/Van Riet; Obied/Ooguri/Spodyneiko/Vafa '18 (see also Bena, Grana, Sethi, Dvali, ....)

• Subsequently, constructions like KKLT and LVS have been subjected to intense scrutiny (with varying success).

Bena/Grana/Van Riet, Van Riet, Moritz/Retolaza/Westphal, Gautason/ Van Hemelryck/Van Riet, Hamada/AH/Shiu/Soler, Bena/Dudas/Grana/Lüst, Lüst/Randall, ... Eventually, serious problems were identified:

- To get a small uplift, throat must be very large (including 'thick')
   Carta/Moritz/Westphal '19
- This leads to strong warping in bulk and 'Singular Bulk Problem'
   Gao/AH/Junghans '20
- Simplest KKLT-models become uncontrolled; Very large tadpole needed for LVS Junghans; Gao/AH/Schrever/Venken '22
- Including effects of curvature at the bottom of the throat makes this much worse

Junghans '22 AH/Schreyer/Venken '22 Venken/Schreyer, Venken ...'24

\_ strong worping singularities in bulk need large throat 1 too much curvature (日)、

Crucial aside:

(Stringy) Quintessence:



 In a nutshell: It does not help! (in spite of many attempts...)

Selection of older and recent work: Cicoli/Pedro/Tasinato/Burgess; Cicoli/DeAlwis/Maharana/Muia/Quevedo; Acharya/Maharana/Muia; Emelin/Tatar; Hardy/Parameswaran; Cicoli/Cunillera/Padilla/Pedro; ....

• One (in my opinion key) argument goes as follows:

(cf. '*F*-Term Problem' in AH/Skrzypek/Wittner '19)

Our world has SUSY broken at TeV, i.e.

 $|F|^2 \sim e^K |DW|^2 \sim \text{TeV}^4$ 

(This part of the superpotential can **not** be rolling to zero – we would see that!)

### What's the status of the 'dS landscape'?:

- Maybe it's there as expected, just in a more complex and less controlled form.... McAllister/Moritz/Nally/Schachner ....
   ...maybe we need an 'F-term-uplift' to replace the anti-D3....
   ...AH/Leonhardt; Krippendorf/Schachner; Lanza/Westphal....
- Maybe string pheno must be completely rethought.
- Very likely: Stringy dS (or even just slow-roll) vacua are more fine-tuned and rare than expected.

The measure problem

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- In any case, the landscape is not likely to become a single vacuum.
  - ⇒ 'Measure problem' (or 'initial condition problem') is still there. Let's revisit it!

Measure problem and potentially decisive role of creation processes

 Standard view: Different vacua → different patches in 'global dS multiverse'. Measure problem ≡ problem of cutoff choice.



Based on the 'Cosmological Central Dogma',

we want to argue for a more Banks '01, Susskind '21 fundamental, quantum-mechanical measure.

Friedrich/AH/Salmhofer/Strauss/Walcher '22, Friedrich/AH/Westphal '24

Towards a 'Quantum-Measure'

• Cosmological Central Dogma:

dS space is a finite system with  $\dim(\mathcal{H}) = e^{S}$ .

• Eternal Inflation  $\equiv$  Series of transitions between different subspaces (with dim $(\mathcal{H}_i) = e^{S_i}$ ).



#### The 'Local Wheeler-DeWitt Measure'

Friedrich/AH/Salmhofer/Strauss/Walcher '22, Friedrich/AH/Westphal '24

- To formalize this 'CCD' perspective, the right approach should be the Wheeler-DeWitt equation.
- Upon gauging time-diffeomorphisms, one has

$$H\psi = i\dot{\psi} \rightarrow H\psi = 0$$

• In our context, the WDW equation needs a source:

$$H\psi = \chi$$

• Such a source term for the creation from nothing is unavoidable since there is also decay to AdS.



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#### The 'Local Wheeler-DeWitt Measure'

- Formally, we have to solve  $H\psi = \chi$  for  $\psi$ and calculate the probability for vacuum dS<sub>i</sub> as  $p_i = ||\psi|_i||^2$ .
- In practice, this reduces to rate equations for a 'flow through the landscape':



(Here any of the 'dS<sub>i</sub>' could also be just a slow-roll plateau)

The outcome is similar to certain 'local measures': Bousso/Freivogel/Yang '06, Garriga/Vilenkin.. '05...'11, Nomura '11, Bousso/Susskind '11, Hartle/Hertog '16

'Local Wheeler-DeWitt Measure' (continued)

- Denote the sources by  $J_i$  and the decay rates by  $\Gamma_{i \rightarrow j}$ .
- Then the relevant rate equations read

$$J_{i} = \sum_{j \in dS} \left( p_{i} \Gamma_{i \rightarrow j} - p_{j} \Gamma_{j \rightarrow i} \right) + p_{i} \sum_{y \in Terminal} \Gamma_{i \rightarrow y} .$$

• The solution can be given as a series:

$$p_i = \frac{1}{\Gamma_i} \left\{ J_i + \sum_j J_j \frac{\Gamma_{j \to i}}{\Gamma_j} + \sum_{j,k} J_j \frac{\Gamma_{j \to k}}{\Gamma_j} \frac{\Gamma_{k \to i}}{\Gamma_k} + \cdots \right\}$$

(Here  $\Gamma_i$  is the total decay rate of vaccum *i*.)

### A conceptual problem: Reheating to Minkowski

- As long as there are only dS and AdS vacua (and a non-zero rate for creation from nothing), finiteness is obvious.
- There is a sensitivity to the number of observers on the horizon-sized patch of the reheating surface.
   But we ignore this (non-exponential) effect.
- However, this changes once we include Minkowski-bubbles: Now we get an infinite reheating surface and no 'natural' cutoff at the horizon:





Our proposal:

• Appeal to an 'Effective CCD', based on the similarity of the reheating surfaces in dS and Minkowski:





• Claim: Even in Minkowski only a finite portion of the surface  $(\sim 1/H_{\rm reh}^3)$  is independent – the rest is gauge redundant.

⇒ Finiteness is regained

## Alternative possibility:

Take infinite Minkowski-space reheating surfaces seriously.
 ⇒ Key prediction: The dark energy in our universe will decay

 our future is Minkowski space.

### Towards explicit predictions

- To solve the rate equations, we obviously need transition rates.
- But, in addition, our local measure crucially depends on creation rates. These depend on End-of-the-World branes:



[Cf. recent discussion of 'Bubble of Something' for String Landscape in Friedrich/AH/Walcher '23. Also, much recent work on inverse 'Bubble of Nothing' process: Garcia-Etxebarria/Montero/Sousa/Valenzuela, Draper et al., Angius/Calderon-Infante/Delgado/Huertas/Uranga, ....]

### ETW Branes

- Traditionally, Hartle-Hawking/Linde-Vilenkin are the leading proposals; alternatives with ETW branes are at best exotic.
- I would argue that the 'Cobordism Conjecture' is one of the most convincing Swampland conjectures; It implies that ETW branes are ubiquitous.

[The conjecture roughly says that the space of geometries is connected. This includes the connection to 'nothing' by a boundary.]

• Rates depend on brane-tensions. Getting those is a challenge...



• Explicitly, creation rates are:  $J \sim exp(\pm S)$  with:



 $\Rightarrow$  For LV sign choice, the 'bos'/'b' creation processes always dominate over 'nb' when the required ETW branes exist.

## Summary / Conclusions

- The problem of realizing (metastable) de Sitter vacua in string theory remains unsolved.
- One is forced to dive into technical details. This is unsatisfcatory. But it's one possibility for making progress ....
- Either way (with or without long-lived dS), predictions need a measure.
- I argued that, in a proper quantum approach, this is sensitive to 'Creation from Nothing'.
- A key 'new' ingredient in this are ETW branes, allowing for 'Bubbles of something' or 'boundary processes'.
- The tensions of those ETW branes are a key research target!